# A Tunable Trisection Bandpass Filter with Constant Fractional Bandwidth Based on Magnetic Coupling

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Abstract — A tunable microstrip trisection bandpass filter with source-load coupling is proposed in this paper. Magnetic cross coupling structure based on varactor loaded open ring resonators is employed. No extra capacitor is used to adjust the coupling coefficients directly. Normalized coupling matrix M is used to calculate the frequency response, and formulas for computing the S parameters based on the resonator with unconventional phase performance are given. Coupling coefficients are investigated by computing the integral of the distributed voltage/current and it is proved that the coupling coefficients can meet the requirements of constant fractional bandwidth (CFBW). Due to the magnetic cross coupled trisection structure and the electrical source-load coupling, this tunable filter has three transmission zeros at finite frequency which can effectively improve the frequency selectivity. The measurement agrees well with the simulation. The center frequency can be tuned from 1097 MHz to 1936 MHz. In this tuning range, the insertion loss varies from 7.8 dB to 4 dB. A constant fractional bandwidth of about 5% is achieved.

*Index Terms* — cross-coupling, trisection filter, tunable bandpass filter.

## I. INTRODUCTION

the fast development of wireless communications, reconfigurable RF/microwave systems, which can be used in electronic countermeasures, software defined radio and so on, are highly demanded. The tunable filter, which dynamically receives the wanted signals or rejects the interferences, is regarded as a core component in reconfigurable RF/microwave systems. For tunable bandpass filters, high frequency selectivity is vital since it can effectively restrain the near-passband interference. One approach to improve frequency selectivity is utilizing high resonators with high unloaded quality factor (Qu) such as cavity resonators [1, 2] or SIW filters [3, 4]. For varactortuning microstrip bandpass filters, the resonator quality factor is relatively lower. Transmission zeros can be used for improving the frequency selectivity, and crossed coupling can bring transmission zeros by making signals transmit through different paths and cancel each other at a specified frequency [5-7].

Crossed coupling has already been widely used in fixed filters [8-12]. By applying source-load coupling, Nth order canonical filters with N finite transmission zeros can be achieved [10]. Recently, tunable crosscoupled tunable filters have been reported. The roll-off rate of second order tunable cross-coupled filter based on source-load cross coupling [13, 14] is limited, and several fourth order cross-coupled filters with higher frequency selectivity have been reported [6, 7, 15, 16]. In [6, 7, 15-17], tunable capacitors are placed in the coupling region for tuning the electric coupling coefficient directly, and bandwidth becomes tunable since coupling coefficients can be tuned. However, this method requires more capacitors, which may lead to higher insertion loss and increase the difficulty of practical layout design. Meanwhile, tunable capacitors which are now available on the market are relatively large as coupling capacitor. Therefore, the capacitance in resonators should be very large to maintain a relatively small coupling coefficient [6]. Tunable third order filter is a useful topology to reduce the designing difficulty and maintain the roll-off rate [18-20]. A typical third order filter is trisection filter which is desirable for some applications requiring only a higher selectivity on one side of the passband due to the advantage of its asymmetrical frequency response. [18-20] reported several third order tunable filters but they only have one or two transmission zeros. Combining cross coupled trisection structure with source-load cross coupling can bring three transmission zeros.

Open ring resonators (ORRs) are common used in microstrip filters [8, 21-23]. It can be utilized for designing tunable filters by adding a tunable capacitor at the open ends of the microstrip line. At the resonant frequency, the magnetic energy is mainly stored in the middle of the microstrip line while the majority of electrical energy is stored at the terminal region of the microstrip line and the capacitor. The larger the capacitance, the less electrical energy there is in the microstrip line, therefore, the electrical coupling can

not make the fractional bandwidth of tunable filter keep in constant.

In this paper, magnetic coupling structure based on varactor loaded ORRs are utilized to design tunable trisection filter with constant fractional bandwidth. Formulas for computing the *S* parameters based on the resonator with unconventional phase performance are given. Coupling coefficients are investigated by computing the integral of the distributed voltage/current. All-magnetic cross coupled resonators and source to load electric coupling are designed to generate three transmission zeros for improving the frequency selectivity.

### II. FILTER SYNTHESIS

The most important parameters that influence filter bandwidth are coupling coefficient k and external quality factor  $Q_{ext}$ . In order to attain the desired k and  $Q_{ext}$  for constant fractional bandwidth, filter synthesis can be applied based on normalized coupling matrix M [8, 24-26]. The formulas for calculating the lowpass S parameters based on M matrix can be written as [8]:

$$\begin{cases} S_{21} = -2jA_{n+2,1}^{-1} \\ S_{11} = 1 + 2jA_{1,1}^{-1} \end{cases}, \tag{1}$$

where *n* is the order of the filter,  $A_{i,j}^{-1}$  denotes the  $i^{th}$  row and  $j^{th}$  column element of  $A^{-1}$ , which is given by:

$$A = M + \Omega \cdot U - j \cdot q, \qquad (2)$$

in which U is the  $(n+2) \times (n+2)$  identity matrix except  $U_{1,1} = U_{n+2,n+2} = 0$ ; q is a  $(n+2) \times (n+2)$  matrix with all zeros except  $q_{1,1} = q_{n+2,n+2} = 1$ , and  $\Omega$  is angular frequency. To get bandpass response, a lowpass-to-bandpass transformation can be written as:

$$\Omega = \frac{1}{FBW} \left( \frac{\omega}{\omega_0} - \frac{\omega_0}{\omega} \right), \tag{3}$$

where  $\omega_0$  is center angular frequency of the transformed bandpass filters,  $\omega$  denotes angular frequency and FBW is the fractional bandwidth. The desired coupling coefficient between  $i^{th}$  and  $j^{th}$  resonators  $k_{ij}$ , and external quality factor  $Q_{ext}$ , which can determine the frequency response of the bandpass filter, can be obtained as:

$$k_{ij} = M_{ij} \cdot FBW = M_{ij} \frac{ABW}{f_0}, \tag{4}$$

$$Q_{ext} = \frac{1}{M_{S1}^2 FBW} = \frac{f_0}{M_{S1}^2 ABW},$$
 (5)

where  $M_{ij}$  denotes the entries in M matrix (i, j = S, 1, 2... N, L, S stands for the input source and L for the output load. The filter is Nth order),  $f_0$  is the center frequency and ABW is the absolute bandwidth.

Figure 1 is the topology used for designing the tunable filter in this paper, where, E stands for electrical coupling and M is magnetic coupling. 1~3 denote resonator 1 to resonator 3. There are three paths from source to load, and the M matrix in (6) is used for

computing the frequency response, as in Fig. 2. This topology can generate three transmission zeros at finite frequency.

$$M = \begin{pmatrix} 0 & 1.3 & 0 & 0 & -0.005 \\ 1.3 & 0 & 0.78 & 0.58 & 0 \\ 0 & 0.78 & 0 & 0.78 & 0 \\ 0 & 0.58 & 0.78 & 0 & 1.3 \\ -0.005 & 0 & 0 & 1.3 & 0 \end{pmatrix}. \tag{6}$$

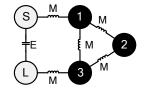


Fig. 1. The topology of the trisection bandpass filter.

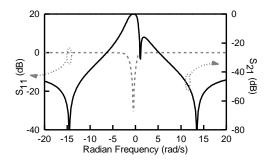


Fig. 2. The frequency response of the *M* matrix.

In the traditional synthesis method, the resonators have a phase shift of  $\pi/2$  at low frequency and  $-\pi/2$  at high frequency [26]. Figure 3 shows that the phase shift of varactor loaded ORR is opposite. Hence the U in (2) should be changed to:

$$U = diag[0, -1, -1, -1, 0]. (7)$$

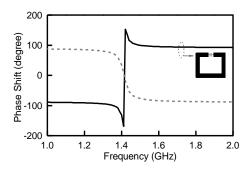


Fig. 3. The phase response of varactor loaded ORR (black solid line), and the resonators in traditional synthesis method (gray dashed line).

Figure 4 shows the response based on the topology shown in Fig. 1 based on the varactor loaded ORR, with

*FBW* = 0.05 under center frequency of 1.3 GHz, 1.5 GHz, and 1.7 GHz, respectively. The desired coupling coefficients and external quality factor versus center frequency can be calculated using (4) and (5), as shown in Fig. 2.

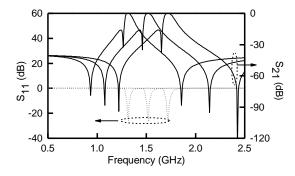


Fig. 1. Bandpass response of the M matrix. This calculation uses the modified U.

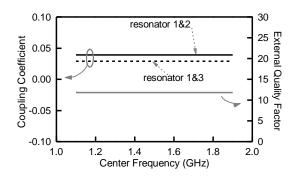


Fig. 2. The desired coupling coefficients and external quality factor extracted from M matrix.

#### III. FILTER DESIGN

Figure 6 shows the tunable cross coupled trisection bandpass filter based on the magnetic coupled varactor loaded ORR.  $D_1 \sim D_6$  are the varactors and  $R_1 \sim R_6$  are 5.1 k $\Omega$  resistors used as RF chokes. Feedlines utilize magnetic coupling to achieve matching and electrical source-load coupling is implemented through feedlines. Table 1 lists their physical parameters in Fig. 6. Three types of geometry parameters are used for studying the coupling structure.

## A. Coupling coefficients

The relationship between center angular frequency  $\omega_0$  and the equivalent capacitance of varactors can be analyzed by half circuit of the resonators. For resonator 1 and resonator 3 in Fig. 6, their half circuit is shown in Fig. 7 (a). It is made up with two short ended microstrip lines and two capacitors named  $C_L$ .  $Y_0$  is the characteristic admittance of the microstrip lines,  $\theta_1$  is

the electrical length of the coupling region and  $\theta_2$  is the electrical length of each microstrip line.

At resonating frequency  $\omega_0$ , the tunable capacitance  $C_L$  can be written as:

$$j\omega C_L - jY_0 \cot \theta_2 = 0, \qquad (8)$$

$$C_L = \frac{Y_0 \cot \theta_2 \big|_{\omega = \omega_0}}{\omega_0} \,. \tag{9}$$

The calculated curves of the tunable capacitor  $C_L$  versus center frequency based on the three geometries in Table 1 are shown in Fig. 8. Since  $L_2$ ,  $L_3$  and  $W_2$  of the three geometries are same, the calculated curves coincide.

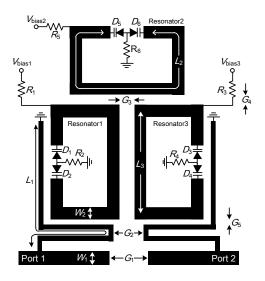


Fig. 6. The electrical circuit model of this filter.

Table 1: Physical parameters of 3rd tunable filters

Parameters	Dimensio	Dimension	Dimensio		
(unit: mm)	n 1	2	n 3		
$L_2$	26.9	26.9	26.9		
$L_3$	11.1	11.1	11.1		
$W_2$	0.7	0.7	0.7		
$G_3$	1.8	2	2.2		
$G_4$	0.4	0.6	0.8		

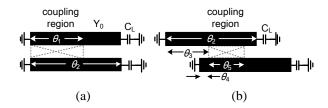


Fig. 7. The equivalent model of half circuit of coupled resonators: (a) resonator 1 and 3, and (b) resonator 1 and 2.

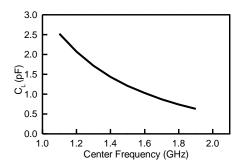


Fig. 8. The calculated  $C_{\rm L}$  versus center frequency.

Coupling coefficients can be analyzed by energy method used in [18, 27]. In the half circuit of resonator 1 and 3 shown in Fig. 7 (a), the electrical energy and magnetic energy stored in each resonator at center frequency can be written as:

$$W_{C_1} = W_{L_1} = W_{C_3} = W_{L_3} = \frac{1}{2} L Y_0^2 \int_0^{\theta_2} \cos^2 \theta d\theta , \qquad (10)$$

where  $W_{Li}$  (i = 1, 3) denotes magnetic energy in  $i^{th}$  resonator,  $W_{Ci}$  (i = 1, 3) denotes electric energy in  $i^{th}$  resonator, L is the self-inductance of the microstrip lines. Meanwhile, the magnetic energy stored in the coupling region,  $W_{Lm}$ , is expressed as:

$$W_{L_m} = L_m Y_0^2 \int_0^{\theta_1} \cos^2 \theta d\theta . \tag{11}$$

 $L_m$  is the mutual inductance between the two lines. Therefore, the magnetic coupling coefficient of this half circuit is obtained by the energy relationship:

$$k_{L} = \frac{W_{L_{m}}}{2\sqrt{W_{L_{1}}W_{L_{1}}}}.$$
 (12)

The electrical energy stored in coupling region is,

$$W_{C_m} = -C_m \int_0^{\theta_1} \sin^2 \theta d\theta . \tag{13}$$

 $C_m$  is the mutual capacitance between the two lines. So electrical coupling coefficient can be calculated:

$$k_C = \frac{W_{C_m}}{2\sqrt{W_{C_1}W_{C_3}}} \quad . \tag{14}$$

The total coupling coefficient of this half circuit is the sum of  $k_C$  and  $k_L$ , which is,

$$k = k_L + k_C (15)$$

As for resonator 1 and resonator 2, the half circuit is shown in Fig. 7 (b). Energy stored in each resonator can be calculated using (10). However, the formula to get energy stored in coupling region should be modified:

$$W_{L_{m}} = L_{m} Y_{0}^{2} \int_{0}^{\theta_{2}} \cos\theta \cos\left(\theta - \left(\theta_{3} - \theta_{4}\right)\right) d\theta, \qquad (16)$$

$$W_{C_m} = -C_m \int_{\theta_1}^{\theta_2} \sin\theta \sin\left(\theta - \left(\theta_3 - \theta_4\right)\right) d\theta . \tag{17}$$

Rewrite (12) and (14) as:

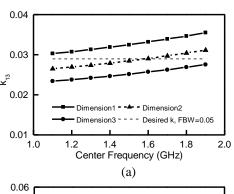
$$k_L = \frac{W_{L_m}}{2\sqrt{W_{L_1}W_{L_2}}},$$
 (18)

$$k_C = \frac{W_{C_m}}{2\sqrt{W_{C_1}W_{C_2}}} \ . \tag{19}$$

For the entire resonators, the energy stored in each resonator is doubled, yet the energy stored in the coupling region does not change. As a result, the coupling coefficient between resonator 1 and 2 is equal to k/2. The coupling coefficients versus center frequency can be calculated by using (9) ~ (19), as shown in Fig. 9 (a) and Fig. 9 (b). It shows that, the strength of coupling can be controlled by adjusting its physical parameters, and the geometry parameters of Dimension 2 in Table 1 will be implemented in the design so that a 5% constant FBW (CFBW) can be achieved.

Simulation is carried out to verify the above calculation.  $f_1$  and  $f_2$  are the splitting resonate frequencies in the simulation, and the coupling coefficient k can be calculated as:

$$k = \frac{f_1^2 - f_2^2}{f_1^2 + f_2^2} \,. \tag{20}$$



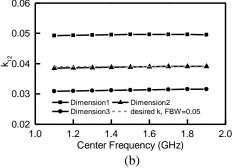


Fig. 9. The calculated results versus center frequency: (a) coupling coefficient  $k_{12}$ , and (b) coupling coefficient  $k_{13}$ .

Figure 10 shows the simulated loaded capacitance versus center frequency compared with calculated results. It shows that the center frequency can be tuned from 1.1 GHz to almost 2 GHz when  $C_L$  is tuned from 0.6 pF to 2.4 pF, and the calculation matches the simulation very well. The slope of coupling coefficient curves is very small, which meets the requirement of

CFBW response.

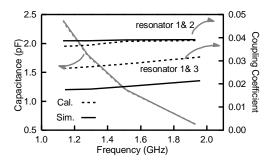


Fig. 10. The tunable capacitance and coupling coefficients versus center frequency of calculation and simulation.

# B. External quality factor

External quality factor  $Q_{ext}$  mainly influences the bandwidth, roll-off, inband ripple. For tunable filters with CFBW,  $Q_{ext}$  needs to be constant. It can be gotten by simulating single loaded circuits shown in Fig. 11 (a) [8], and  $Q_{ext}$  can be calculated as:

$$Q_{ext} = \frac{2\pi f_0 \tau_{S11}}{4} \,, \tag{21}$$

or,

$$Q_{ext} = \frac{f_0}{\Delta f_{\pm \pi/20}} = \frac{f_0}{f_{-\pi/2} - f_{+\pi/2}},$$
(22)

where  $\tau_{s11}$  is group delay of  $S_{11}$ ,  $f_{\pi/2}$  and  $f_{-\pi/2}$  are frequency point at which the phase of  $S_{11}$  of the one-terminal network in Fig. 11 (a) is  $\pm \pi/2$ , respectively.

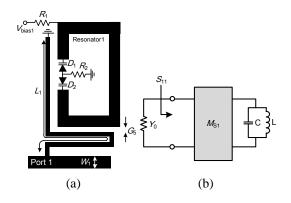


Fig. 11. The simulation circuit and its equivalent model.

 $L_1$ =25 mm,  $W_1$ =2.2 mm and other parameters except  $G_5$  are kept as in dimension 2 mentioned earlier. The gap between feedline and resonator can influence  $Q_{ext}$ . We simulated four conditions that  $G_5$  in Fig. 11 (a) is 0.05 mm, 0.1 mm, 0.2 mm and 0.3 mm, respectively. The simulated external quality factors are shown in Fig. 12. When  $G_5$  increases,  $Q_{ext}$  decreases. This can be explained by the equivalent model in Fig. 11 (b). At  $f_{-\pi/2}$  or  $f_{\pi/2}$ , we have:

$$Y_0 = \left| j \frac{2\pi f_{\pm\pi/2} L M_{S1}^2}{1 - (2\pi f_{\pm\pi/2})^2 L C} \right|, \tag{23}$$

because,

$$S_{11} = \frac{Y_0 - j \frac{\omega L M_{S1}^2}{1 - \omega^2 L C}}{Y_0 + j \frac{\omega L M_{S1}^2}{1 - \omega^2 L C}}$$
(24)

By solving equation (23), the denominator of equation (22) is attained:

$$f_{-\pi/2} - f_{+\pi/2} = \frac{M_{S1}^2}{2\pi C Y_0} \ . \tag{25}$$

Substituting (25) into (22), the external quality of the circuits in Fig. 11 (b) is:

$$Q_{ext} = \frac{CY_0}{M_{SI}^2 \sqrt{LC}} \quad . \tag{26}$$

when the gap becomes narrower, the coupling is enhanced so  $|M_{S1}|$  increases. Equation (26) implies that the external quality factor will decrease if C and L are constant (i.e.,  $f_0$  is fixed). When  $G_5$  equals to 0.1 mm, the external quality factor is the closest to the desired value.  $Q_{ext}$  is also required to be unchangeable for CFBW filters and the simulated results satisfy this requirement.

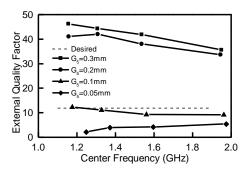


Fig. 12. Simulated external quality factor versus center frequency.

#### C. Source-load coupling

 $S_{21}$  of lowpass prototype filter at  $\Omega = \infty$  can be written as [26]:

$$S_{21\infty} = \frac{2|M_{SL}|}{1 + M_{SL}^2} \ . \tag{27}$$

If the coupling of source-load is enhanced,  $|M_{SL}|$  will increase, and  $S_{21\infty}$  increases when  $|M_{SL}| < 1$ . The electric type  $M_{SL}$  of this filter is desired to be -0.005 for a 40-dB rejection level.  $M_{SL}$  can be simulated by removing all resonators. A two-port network consisting of only two feed lines is obtained as shown in Fig. 13. The image part of  $Y_{12}$  of the network is used to extract  $M_{SL}$  [28] since  $Y_{12}$  denotes a J inverter with J=Im( $Y_{12}$ ):

$$M_{SL} = \text{Im}(Y_{12}) \times 50$$
. (28)

The simulation is carried out when  $G_1$  (other parameters in Fig. 6 are fixed,  $G_2$ =2.4 mm) is varying. Figure 14 shows the different  $Im(Y_{12})$ . It is obvious that  $|Im(Y_{12})|$  is negatively relative to  $G_1$ .  $|Im(Y_{12})|$  is controllable so it is possible to meet desired  $|M_{SL}|$  by adjusting the gap. Finally,  $G_1$  is chosen to be 3.4 mm.

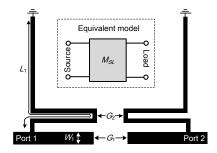


Fig. 13. The simulation circuit of  $M_{\rm SL}$ .

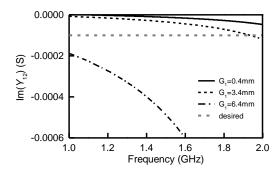


Fig. 14. The simulated imaginary of  $Y_{12}$ .

# IV. FILTER SIMULATION

A layout of the filter is designed and simulated. The substrate with h = 0.8 mm,  $\varepsilon_r = 2.65$ ,  $\tan \delta = 0.001$ , and SMV1405 varactor diode are utilized. Figure 15 shows the *S* parameters under 0V, 1V, 3V, 7V, 14V and 30V bias voltages.

The center frequency varies from 1086MHz to 1941 MHz which means a 179% tuning range is achieved. As shown in Fig. 16, in the tuning range, fractional bandwidth of this filter keeps almost unchangeable (4.25% ~4.73%) and the insertion loss is less than 6.2 dB. Additionally, there are three transmission zeros besides the passband. One is generated by the utilized trisection structure and is located at lower edge of the passband. The other two transmission zeros exist at each edge of the passband due to the source-load coupling. These zeros can improve the selectivity of this filter. The rejection is approximately 40 dB (higher frequency, estimated by the frequency response's asymptote in Fig. 15) or higher than 50 dB (lower frequency, estimated by the frequency response's stationary point in Fig. 15).

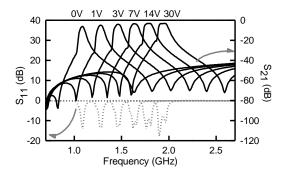


Fig. 15. Simulated S parameters under typical tuning states.

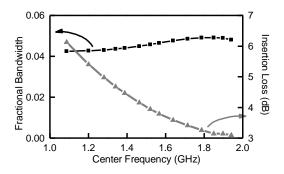


Fig. 16. Simulated insertion loss and fractional bandwidth (FBW).

# V. FABRICATION AND MEASUREMENT

The filter is fabricated with an overall size of  $36 \text{ mm} \times 40 \text{ mm}$ , as shown in Fig. 17. Three lines on the sectorial pads are connected to the DC supply to bias the varactors. The frequency response of the filter is measured by Agilent E5071C vector network analyzer.

The bias voltage is swept from 0 V to 30 V and  $V_{\rm bias1} = V_{\rm bias2} = V_{\rm bias3}$ . Figure 18 shows the measured S parameters with 0V, 1V, 3V, 7V, 14V, and 30V bias voltages. The measured frequency responses agree well with the simulation. The tuning ratio of center frequency is approximately 176% (1097 MHz  $\sim$  1936 MHz). Three transmission zeros are generated, and the selectivity of this filter is improved by the transmission zeros. The rejection level is about 40 dB (higher frequency) or higher than 50 dB (lower frequency).

Figure 19 shows the measured insertion loss and fractional bandwidth. The fractional bandwidth maintains 4.1% ~5.3%. It is very narrow, which is useful for channelization. The insertion loss changes from 7.8 dB to 4 dB, which is approximately 1 dB higher than simulation. Besides the inaccuracy in PCB fabrication and the practical discrepancy between varactors, this is mainly due to the errors in welding. The PCB is relatively small in size and the gap between feedlines and resonators is very narrow. So, the solder joint,

which is not small enough to be ignored or to be regarded as ideal connection point, has a great effect on the performance. A similar situation can be seen in [29]. This is not simulated in software and can be solved by precise fabrication or lower frequency application with larger board area. Nonlinear characteristics is usually represented by IIP3 (input  $3^{rd}$  order intercept point), for example [29]. It can be obtained by feeding dual tone signals and measuring the output frequency spectrum. The frequencies of the input signals are set to be  $f_0 \pm 1$  MHz. The tested IIP 3 is shown in Fig. 20. Table 2 shows the comparison with several recent cross-coupled tunable filters.

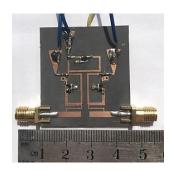


Fig. 17. The photograph of the fabricated filter.

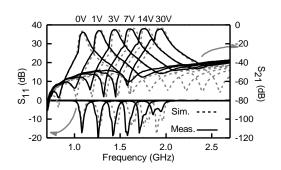


Fig. 18. The measured S parameters.

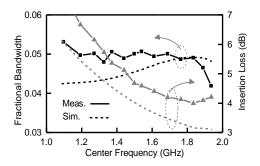


Fig. 19. The measured insertion loss and FBW of this filter.

Table 2: Com	parison	with	recent	cross-	-couple	ed tuna	able f	ilters

Reference	Tuning Range (GHz)	Insertion Loss (dB)	Bandwidth	Number of Zeros
[13]	1.11~1.51	3.6~4.2	29±3MHz*, 1.8%~2.6%*	2
[14]	1.15~2	3.6~2.4 114±4MHz*, 9.9%~5.7%*		2
[15]	1.5~1.75	4.5~11	4.5~11 variable	
[16]	1.25~2.1	3.5~8.5@FBW=4% 3.5~6.5@FBW=5.5%	variable	2
[19]	0.41~0.82	<6.5	9±1%	1
[30]	0.95~1.46	1.8~3.6	12.5±2%	3
[31]	1.36~1.78	2.8~5	93±7 MHz	2
This work	1.097~1.936	7.8~4	4.7±0.6%	3

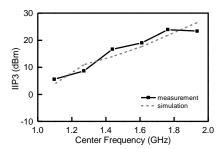


Fig. 20. The measured IIP3 compared with simulation.

It should be noted that the default bandwidth is 3dB bandwidth while the marked data using the symbol \* is 1 dB bandwidth in Table 2. Different from the 3<sup>rd</sup> order filter in [30] with 3 transmission zeros, this

crossed coupled filter does not require additional coupling coefficient tuning elements. This is also an advantage when competing with 4<sup>th</sup> order filter in [15, 16]. Filter in [31] do not need additional capacitors, either, but it needs two different bias voltage values while in this filter only need one bias voltage value.

# VI. CONCLUSION

In this paper, a cross coupled tunable trisection filter with source-load coupling is proposed. All magnetic coupling is employed in trisection structure while source and load are electrically coupled. Three transmission zeros that can effectively improve the frequency selectivity are generated due to the crossed coupled topology. *M* matrix is used to guide this design. The phase performance of the resonator is studied then

the computing formula of *S* parameters based on *M* is modified for this resonator. Energy relationship is utilized to analyze the coupling coefficients, which proves that the designed structures are able to satisfy the requirements of CFBW. According the theoretical analysis, coupling coefficients can be controlled and meet the desired values. Source-load coupling, and external quality factor are studied, optimized by simulation. The filter is fabricated on F4B-2 substrate and measured by a vector network analyzer. The measurement matches the simulation well. The tuning range is 1097 MHz to 1936 MHz, meaning that a 176% tuning ratio is achieved. The fractional bandwidth is keeping about 5%. In tuning range, the insertion loss varies from 7.8 dB to 4 dB.

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