# Accurate Far-Field Measurement of Wide-Angle Scanning Phased Array Antennas 

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#### Abstract

This paper presents a new approach to the accurate determination of the far-field (FF) pattern of wide-angle scanning phased array antennas (PAAs) that uses planar near-field (NF) technology. When PAAs are engaged in wide-angle scanning, a finite scanning plane and inaccurate probe compensation can lead to truncation errors. The proposed approach reduces these by rotating the PAA to make its main beam perpendicular to the scanning plane. To verify the effectiveness of the approach, both theoretical calculations and actual measurements were undertaken. The simulated and measured results prove that the proposed method is more accurate than traditional NF technology for a large PAA scanning angle, such as $\psi_{0}=50^{\circ}$ and $\Delta_{0}=50^{\circ}$ (in simulation; $54^{\circ}$ when measured).


Index Terms - Antenna measurement, planar near-field technology, wide-angle scanning phased array antenna.

## I. INTRODUCTION

The effective measurement of antenna technologies [1-4] has always been a key issue. Phased array antennas (PAAs) are antennas that can be electrically controlled to change the direction and shape of the beam. In addition to fifth generation (5G) millimeter-wave mobile terminals [5], PAAs are often required in modern radar applications to provide scanning with spatial coverage over a broader angular range. A prominent concern in the design of PAAs has been to enlarge their scanning angle, so measuring their wide-angle scanning performance has become an important topic of research [6]. Wide-angle scanning PAA measurement technologies are usually applied to either far-field (FF) measurements or nearfield (NF) measurements [7-9]. However, individual FF measurements tend to deliver very limited information. NF approaches are therefore commonly used for antenna radiation and RCS measurements. Existing planar NF measurements are also best suited to measuring antennas with pencil-beam characteristics. However, finite scanning surface truncation errors often affect the
accuracy of planar NF measurements. For antennas that do not scan the beam and antennas that do not have a large scan angle, there usually needs to be a cutoff at the edge of the planar NF scanning plane, below -40 dB , to minimize the truncation error. When the beam of an antenna under test (AUT) is scanned, a large truncation angle is often required to meet the -40 dB cutoff. Unfortunately, large truncation angles are not always possible during actual measurement processes. Furthermore, when the scanning angle of an AUT is large, there can also be inaccuracies in the calculation of the probe compensation.

Various efforts have been made to try and reduce the truncation errors. The most commonly-used method is to select the most appropriate window function. This will have a smooth taper at the edges of the scan plane [10]. However, this method reduces the extent of the reliable region. An equivalent electric current or magnetic current can be introduced to determine the FF pattern [11], but this method requires certain kinds of a priori information, such as the geometry and position of the antenna. It is also extremely time-consuming for electrically large antennas. Some previously proposed techniques [e.g., 12] have suggested enlarging the reliable angular region, but they increase the complexity of the scanning system by demanding additional measurements of the shape. Recently, band-limited signal extrapolation algorithms have been applied to radiation [13] and RCS measurements [14] in the NF region, so as to reduce the truncation errors. However, these approaches cannot be directly applied to the measurement of wide-angle scanning antennas because most of the energy remains outside the NF scanning plane of the AUT and cannot be ignored.

In this paper, we present a novel method for accurately measuring wide-angle scanning PAAs. Prior to the planar NF measurement, the AUT is rotated by a specific angle so that the main beam is perpendicular to the scanning plane. During NF scanning on the test plane, it is then relatively easy to obtain lower near-field
truncation levels. After the mathematical transformation introduced in this paper, it is possible to achieve accurate PAA FF measurements.


Fig. 1. The relationship between $0-x y z$ and $o-x^{\prime} y^{\prime} z^{\prime}$.

## II. THEORETICAL BACKGROUND

## A. Coordinate system transformation

Figure 1 illustrates the relationship between an original coordinate system, o-xyz, and a translated coordinate system, o-x'y'z'. Assuming that the rotation angles are $\psi_{0}$ and $\Delta_{0}$ in the original coordinate system, we can get the $o-x^{\prime} y^{\prime} z^{\prime}$ coordinate system after two rotations. This relationship can be written as:

$$
\begin{gather*}
{\left[\begin{array}{l}
\hat{x}^{\prime} \\
\hat{y}^{\prime} \\
\hat{z}^{\prime}
\end{array}\right]=T\left[\begin{array}{l}
\hat{x} \\
\hat{y} \\
\hat{z}
\end{array}\right]=\left[\begin{array}{lll}
t_{11} & t_{12} & t_{13} \\
t_{21} & t_{22} & t_{23} \\
t_{31} & t_{32} & t_{33}
\end{array}\right]\left[\begin{array}{l}
\hat{x} \\
\hat{y} \\
\hat{z}
\end{array}\right],}  \tag{1a}\\
{\left[\begin{array}{l}
\hat{x} \\
\hat{y} \\
\hat{z}
\end{array}\right]=\left[\begin{array}{lll}
t_{11} & t_{21} & t_{31} \\
t_{12} & t_{22} & t_{32} \\
t_{13} & t_{23} & t_{33}
\end{array}\right]\left[\begin{array}{l}
\hat{x}^{\prime} \\
\hat{y}^{\prime} \\
\hat{z}^{\prime}
\end{array}\right],} \tag{1b}
\end{gather*}
$$

where

$$
t_{11}=\cos \psi_{0}, t_{12}=0, t_{13}=-\sin \psi_{0},
$$

$t_{21}=-\sin \Delta_{0} \sin \psi_{0}, t_{22}=\cos \Delta_{0}, t_{23}-\sin \Delta_{0} \cos \psi_{0}$
$t_{31}=\cos \Delta_{0} \sin \psi_{0}, t_{32}=\sin \Delta_{0}, t_{33}=\cos \Delta_{0} \cos \psi_{0}$. (2)
Assuming that the coordinates of the vector, $\vec{r}$, in the two coordinate systems are $(x, y, z)$ and $\left(x^{\prime}, y^{\prime}, z^{\prime}\right)$, it can be represented as follows:

$$
\begin{align*}
\vec{r}= & {\left[\begin{array}{lll}
x & y & z
\end{array}\right]\left[\begin{array}{l}
\hat{x} \\
\hat{y} \\
\hat{z}
\end{array}\right]=\left[\begin{array}{lll}
x^{\prime} & y^{\prime} & z^{\prime}
\end{array}\right]\left[\begin{array}{l}
\hat{x}^{\prime} \\
\hat{y}^{\prime} \\
\hat{z}^{\prime}
\end{array}\right]=} \\
& {\left[\begin{array}{lll}
x^{\prime} & y^{\prime} & z^{\prime}
\end{array}\right]\left[\begin{array}{lll}
t_{11} & t_{12} & t_{13} \\
t_{21} & t_{22} & t_{23} \\
t_{31} & t_{32} & t_{33}
\end{array}\right]\left[\begin{array}{l}
\hat{x} \\
\hat{y} \\
\hat{z}
\end{array}\right] . } \tag{3}
\end{align*}
$$

The relationship between $(x, y, z)$ and $\left(x^{\prime}, y^{\prime}, z^{\prime}\right)$ will then be:

$$
\left[\begin{array}{l}
x  \tag{4}\\
y \\
z
\end{array}\right]=\left[\begin{array}{lll}
t_{11} & t_{21} & t_{31} \\
t_{12} & t_{22} & t_{32} \\
t_{13} & t_{23} & t_{33}
\end{array}\right]\left[\begin{array}{l}
x^{\prime} \\
y^{\prime} \\
z^{\prime}
\end{array}\right]=\left[\begin{array}{l}
t_{11} x^{\prime}+t_{21} y^{\prime}+t_{31} z^{\prime} \\
t_{12} x^{\prime}+t_{22} y^{\prime}+t_{32} z^{\prime} \\
t_{13} x^{\prime}+t_{23} y^{\prime}+t_{33} z^{\prime}
\end{array}\right] .
$$

Similarly, we can obtain:

$$
\left[\begin{array}{l}
x^{\prime}  \tag{5}\\
y^{\prime} \\
z^{\prime}
\end{array}\right]=\left[\begin{array}{lll}
t_{11} & t_{12} & t_{13} \\
t_{21} & t_{22} & t_{23} \\
t_{31} & t_{32} & t_{33}
\end{array}\right]\left[\begin{array}{l}
x \\
y \\
Z
\end{array}\right]=\left[\begin{array}{l}
t_{11} x+t_{12} y+t_{13} z \\
t_{21} x+t_{22} y+t_{23} z \\
t_{31} x+t_{32} y+t_{33} z
\end{array}\right] .
$$

Vectors in any direction in space $(\psi, \Delta)$ will correspond to the $o-x y z$ coordinate system and ( $\psi^{\prime}, \Delta^{\prime}$ ) will correspond to the $o-x^{\prime} y^{\prime} z^{\prime}$ coordinate system. Any unit vector, $\hat{r}$, can be represented as follows in the two coordinate systems:

$$
\begin{array}{r}
\hat{r}=\left[\begin{array}{lll}
\cos \Delta \sin \psi & \sin \Delta & \cos \Delta \cos \psi
\end{array}\right]\left[\begin{array}{l}
\hat{x} \\
\hat{y} \\
\hat{z}
\end{array}\right] \\
\hat{r}=\left[\begin{array}{lll}
\cos \Delta^{\prime} \sin \psi^{\prime} & \sin \Delta^{\prime} & \cos \Delta^{\prime} \cos \psi^{\prime}
\end{array}\right]\left[\begin{array}{l}
\hat{x}^{\prime} \\
\hat{y}^{\prime} \\
\hat{z}^{\prime}
\end{array}\right] \\
=\left[\begin{array}{lll}
\cos \Delta^{\prime} \sin \psi^{\prime} & \sin \Delta^{\prime} & \cos \Delta^{\prime} \cos \psi^{\prime}
\end{array}\right] T\left[\begin{array}{l}
\hat{x} \\
\hat{y} \\
\hat{z}
\end{array}\right] . \tag{6}
\end{array}
$$

So, the transformation relationship between $\left(\psi^{\prime}, \Delta^{\prime}\right)$ and $(\psi, \Delta)$ will be:

$$
\Delta=\arcsin \left(t_{22} \sin \Delta^{\prime}+t_{32} \cos \Delta^{\prime} \cos \psi^{\prime}\right)
$$

$\psi=\operatorname{angle}\left(\left(t_{13} \cos \Delta^{\prime} \sin \psi^{\prime}+t_{23} \sin \Delta^{\prime} t_{33} \cos \Delta^{\prime} \cos \psi^{\prime}\right)\right.$
$\left.+j\left(t_{11} \cos \Delta^{\prime} \sin \psi^{\prime}+t_{21} \sin \Delta^{\prime}+t_{31} \cos \Delta^{\prime} \cos \psi^{\prime}\right)\right)$. (7)
Similarly, the transformation relationship between $(\psi, \Delta)$ and $\left(\psi^{\prime}, \Delta^{\prime}\right)$ will be:

$$
\begin{gather*}
\Delta^{\prime}=\arcsin \left(t_{21} \cos \Delta \sin \psi+t_{22} \sin \Delta+t_{23} \cos \Delta \cos \psi\right) \\
\psi^{\prime}=\operatorname{angle}\left(\left(t_{31} \cos \Delta \sin \psi+t_{32} \sin \Delta+t_{33} \cos \Delta \cos \psi\right)\right. \\
\left.\quad+j\left(t_{11} \cos \Delta \sin \psi+t_{13} \cos \Delta \cos \psi\right)\right) . \tag{8}
\end{gather*}
$$

The relationship between $(\hat{\psi}, \hat{\Delta})$ and $(\hat{x}, \hat{y}, \hat{z})$ is thus:

$$
\left[\begin{array}{c}
\hat{\psi}  \tag{9}\\
\hat{\Delta}
\end{array}\right]=\left[\begin{array}{ccc}
\cos \psi & 0 & -\sin \psi \\
-\sin \Delta \sin \psi & \cos \Delta & -\sin \Delta \cos \psi
\end{array}\right]\left[\begin{array}{l}
\hat{x} \\
\hat{y} \\
\hat{z}
\end{array}\right]
$$

Similarly,

$$
\left[\begin{array}{c}
\hat{\psi}^{\prime}  \tag{10}\\
\hat{\Delta}^{\prime}
\end{array}\right]=\left[\begin{array}{ccc}
\cos \psi^{\prime} & 0 & -\sin \psi^{\prime} \\
-\sin \Delta^{\prime} \sin \psi^{\prime} & \cos \Delta^{\prime} & -\sin \Delta^{\prime} \cos \psi^{\prime}
\end{array}\right]\left[\begin{array}{c}
\hat{x}^{\prime} \\
\hat{y}^{\prime} \\
\hat{z}^{\prime}
\end{array}\right],
$$

The dot matrix between $(\hat{\psi}, \hat{\Delta})$ and $\left(\hat{\psi}^{\prime}, \widehat{\Delta}^{\prime}\right)$ can be expressed as:

$$
\hat{\psi} \cdot \hat{\psi}^{\prime}=\cos \psi^{\prime}\left(\mathrm{t}_{11} \cos \psi-\mathrm{t}_{13} \sin \psi\right) \sin \psi^{\prime}\left(\mathrm{t}_{31} \cos \psi-\right.
$$

$$
\begin{gather*}
\hat{\psi} \cdot \hat{\Delta}^{\prime}=\cos \Delta^{\prime}\left(\mathrm{t}_{21} \cos \psi-\mathrm{t}_{23} \sin \psi\right)-\sin \Delta^{\prime} \cos \psi^{\prime} \\
\left(\mathrm{t}_{31} \cos \psi-\mathrm{t}_{33} \sin \psi\right)-\sin \Delta^{\prime} \sin \psi^{\prime}\left(\mathrm{t}_{11} \cos -\mathrm{t}_{13} \sin \psi\right) \tag{11b}
\end{gather*}
$$

$\widehat{\Delta} \cdot \hat{\psi}^{\prime}=\sin \psi^{\prime}\left(\mathrm{t}_{31} \sin \Delta \sin \psi-\mathrm{t}_{32} \cos \Delta+\right.$
$\left.\mathrm{t}_{33} \sin \Delta \cos \psi\right)-\cos \psi^{\prime} \sin \Delta\left(\mathrm{t}_{11} \sin \psi+\mathrm{t}_{13} \cos \psi\right)$,
$\widehat{\Delta} \cdot \hat{\Delta}^{\prime}=\sin \Delta^{\prime} \cos \psi^{\prime}\left(t_{31} \sin \Delta \sin \psi-t_{32} \cos \Delta+\right.$
$\left.t_{33} \sin \Delta \cos \psi\right)-\cos \Delta^{\prime}\left(t_{21} \sin \Delta \sin \psi-t_{22} \cos \Delta+\right.$
$\left.t_{23} \sin \Delta \cos \psi\right)+\sin \Delta^{\prime} \sin \psi^{\prime} \sin \Delta\left(t_{11} \sin \psi+t_{13} \cos \psi\right)$.

## B. Pattern transformation

If a PAA is measured in an $o-x^{\prime} y^{\prime} z^{\prime}$ coordinate system, then it must be transformed to an $o-x y z$ coordinate system to obtain its characteristics. The corresponding pattern in the two coordinate systems can be expressed by $(\hat{\psi}, \hat{\Delta})$ and $\left(\hat{\psi}^{\prime}, \hat{\Delta}^{\prime}\right)$, respectively:

$$
\begin{gather*}
\vec{F}=F_{\psi^{\prime}}\left(\psi^{\prime}, \Delta^{\prime}\right) \hat{\psi}^{\prime}+F_{\Delta^{\prime}}\left(\psi^{\prime}, \Delta^{\prime}\right) \hat{\Delta}^{\prime}  \tag{12a}\\
\vec{F}=F_{\psi}(\psi, \Delta) \hat{\psi}+F_{\Delta}(\psi, \Delta) \hat{\Delta} . \tag{12b}
\end{gather*}
$$

From this, the transformation relationship between the
$o-x^{\prime} y^{\prime} z^{\prime}$ coordinate system's pattern and the $o-x y z$ coordinate system's pattern can be obtained:

$$
F_{\psi}(\psi, \Delta)=F_{\psi^{\prime}}\left(\psi^{\prime}, \Delta^{\prime}\right)\left(\hat{\psi}^{\prime} \cdot \hat{\psi}\right)+F_{\Delta^{\prime}}\left(\psi^{\prime}, \Delta^{\prime}\right)\left(\hat{\Delta}^{\prime} \cdot \hat{\psi}\right)
$$

$$
\begin{equation*}
F_{\Delta}(\psi, \Delta)=F_{\psi^{\prime}}\left(\psi^{\prime}, \Delta^{\prime}\right)\left(\hat{\psi}^{\prime} \cdot \widehat{\Delta}\right)+F_{\Delta^{\prime}}\left(\psi^{\prime}, \Delta^{\prime}\right)\left(\hat{\Delta}^{\prime} \cdot \widehat{\Delta}\right) \tag{13a}
\end{equation*}
$$

In Eqs. (13a) and (13b), $F_{\psi^{\prime}}\left(\psi^{\prime}, \Delta^{\prime}\right)$ and $F_{\Delta^{\prime}}\left(\psi^{\prime}, \Delta^{\prime}\right)$ can be obtained through measurement, then, by substituting Eqs. (11a-11d) into Eqs. (13a) and (13b), the FF of the PAA can be acquired.

## III. SIMULATED AND MEASURED RESULTS

## A. Calculated results

In order to assess the effectiveness of the proposed approach, a measurement model was first established. Fig. 2 shows a schematic diagram of the model. The AUT here consisted of $N_{x} \times N_{y}$ electric dipole elements. The element intervals were $d_{x}$ and $d_{y}$, respectively, and the amplitude had a -25 dB sidelobe Taylor distribution in the direction of the azimuth and a -35 dB sidelobe Taylor distribution in the direction of the elevation. The distance between the scan plane and the AUT was d. The reliable angles were $\theta_{x}$ and $\theta_{y}$, respectively, and the sampling intervals were $\Delta_{x}$ and $\Delta_{y}$. The details of the parameters are presented in Table 1.


Fig. 2. Schematic diagram of the model.
Table 1: Detailed parameters

| Parameter | Value | Parameter | Value |
| :---: | :---: | :---: | :---: |
| $N_{x}$ | 20 | $\theta_{x}$ | $60^{\circ}$ |
| $N_{y}$ | 20 | $\theta_{y}$ | $60^{\circ}$ |
| $d_{x}$ | $0.6 \lambda$ | $\Delta x$ | $0.4 \lambda$ |
| $d_{y}$ | $0.6 \lambda$ | $\theta y$ | $0.4 \lambda$ |
| $d$ | $5 \lambda$ | Antenna Size <br> Scan Plane Size | $12 \lambda \times 12 \lambda .3 \times 29.3 \lambda$ |

It was assumed that the antenna beam scanning angle was $\psi_{0}=50^{\circ}$ and $\Delta_{0}=50^{\circ}$, with truncation angles of $60^{\circ}$ and $85^{\circ}$, respectively. A comparison of these two cases relating to a normal NF-FF transformation and the theoretical results is shown in Fig. 3. The results show that when the PAA was in a wide-angle scanning mode, the difference between the calculated FF pattern and the theoretical FF pattern decreased as the truncation angle increased. When the truncation angle was $85^{\circ}$, there was an obvious difference between the calculated result
and the theoretical result. However, this truncation angle is impossible to realize in actual practice. When the antenna aperture was rotated by $\psi=50^{\circ}$ and $\Delta=50^{\circ}$, which was the same as the scanning angle of the PAA, the main beam was orthogonal to the scanning plane. Here, a $60^{\circ}$ truncation angle was assumed. As can be seen in Fig. 3 (a), the traditional $60^{\circ}$ truncated angle NF-FF was noticeably different from the theoretical FF. Figure 3 (b) shows that they were in good agreement when the truncated angle increased to $85^{\circ}$. However, as can be seen in Fig. 3 (c), the results relating to the proposed method were consistent with the theoretical results when the truncated angle was just $60^{\circ}$. Figure 4 also shows that the calculated results using the proposed method were in good agreement with the theoretical results. This proves the validity and accuracy of the proposed method.


Fig. 3. Comparison between the theoretical results and calculated results for the elevation plane when $\psi_{0}=50^{\circ}$. (a) Comparison between the theoretical FF and a $60^{\circ}$ truncated angle NF-FF. (b) Comparison between the theoretical FF and an $85^{\circ}$ truncated angle NF-FF. (c) Comparison between the theoretical FF and the proposed method.


Fig. 4. Comparison between the theoretical results and calculated results for the elevation plane when $\Delta_{0}=50^{\circ}$. (a) Comparison between the theoretical FF and a $60^{\circ}$ truncated angle NF-FF. (b) Comparison between the theoretical FF and an $85^{\circ}$ truncated angle NF-FF. (c) Comparison between the theoretical FF and the proposed method.

## B. Measured results

The proposed method was also used to measure a helix antenna array at 2.2 GHz . This was composed of $2 \times 8$ helix antenna elements spaced 85 mm apart, with a uniform distribution. It was fed by a one-to-sixteen power divider. The size of the antenna array was $900 \mathrm{~mm} \times 230 \mathrm{~mm} \times 110 \mathrm{~mm}$, as shown in Fig. 5 (a). Beam scanning at $54^{\circ}$ was realized by changing the cable length between the power divider and the radiation element (Fig. 5 (b)). The antenna was rotated so that the main beam was perpendicular to the scanning surface for another planar NF measurement, as shown in Fig. 5 (c). The positional accuracy of the scanning system was 0.05 mm and its RMS value was 0.1 mm . The $54^{\circ}$
scanning angle results shown in Fig. 6 again confirm the effectiveness of the proposed algorithm.

The simulation-based and measured results together serve to confirm that the proposed method is more accurate than using traditional NF technology when the scanning angle of a PAA is large.


Fig. 5. The antenna and measurement process: (a) the antenna, (b) conducting the FF measurement, and (c) conducting the NF measurement.


Fig. 6 Comparison of the NF-FF results and FF results for a $54^{\circ}$ scanning angle: (a) using the traditional NF-FF method, and (b) using the proposed method.

## IV. CONCLUSION

This paper has presented both theoretical analysis and the real application of an approach that can accurately determine the FF pattern of PAAs based on planar NF technology. Using this approach, it is only necessary to rotate the PAAs, so that the main beam is perpendicular to the scanning plane. The results from both simulations and physical measurements have established the validity of the proposed algorithm. The approach presented in this paper is especially suitable for wide-angle scanning PAAs.

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